Cosmological interpretation for the stochastic signal in pulsar timing arrays

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Based on 2307.00722 (SCPMA), 2307.03141 (SCPMA), 2312.01824 (PRD)

Hunan Normal University

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Cosmological sources

Outline



SMBHB

- Cosmological sources
 - Phase transition
 - Cosmic string
 - Scalar-induced GW



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The Nobel Prize in Physics 2017







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- New era of GW astronomy
- Multi-messenger astronomy



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Pulsar and pulsa	r timing array (PT	A)	
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- Pulsars are highly magnetized, rotating neutron stars that emit regular pulses of electromagnetic radiation.
- GWs can cause tiny distortion in spacetime inducing variations in the time of arrivals (ToAs).
- A PTA pursues to detect nHz GWs by regularly monitoring ToAs from an array of the ultra rotational stable millisecond pulsars.

Introduction

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Timing residual induced by a GWB





$$\begin{aligned} z(t,\hat{\Omega}) &= \frac{\nu_e - \nu_p}{\nu_p} \\ &= \frac{\hat{p}^i \hat{p}^j}{2(1 + \hat{\Omega} \cdot \hat{p})} \left[h_{ij} \left(t_p, \hat{\Omega} \right) - h_{ij} \left(t_e, \hat{\Omega} \right) \right] \\ z(t) &= \int_{S^2} d\hat{\Omega} \, z(t, \hat{\Omega}) \end{aligned}$$

• Timing residual in frequency-domain

$$\tilde{r}(f,\hat{\Omega}) = \frac{1}{2\pi i f} \left(1 - e^{-2\pi i f L(1+\hat{\Omega}\cdot\hat{p})} \right) \times \sum_{A} h_A(f,\hat{\Omega}) F^A(\hat{\Omega})$$

Antenna pattern

$$F^{A}(\hat{\Omega}) = e^{A}_{ij}(\hat{\Omega}) \frac{\hat{p}^{i}\hat{p}^{j}}{2(1+\hat{\Omega}\cdot\hat{p})}$$

Detecting a GWB with PTA

Introduction

• Assume the GWB is isotropic, unpolarized, and stationary

$$\left\langle h_A^*(f,\hat{\Omega})h_{A'}(f',\hat{\Omega}')\right\rangle = \frac{3H_0^2}{32\pi^3 f^3} \delta^2(\hat{\Omega},\hat{\Omega}')\delta_{AA'}\delta(f-f')\Omega_{\rm gw}(f)$$

Spectrum of GWB

$$\Omega_{\rm gw}(f) \equiv \frac{1}{\rho_{\rm crit}} \frac{d\rho_{\rm gw}}{d\ln f}, \qquad \rho_{\rm crit} \ = \frac{3H_0^2}{8\pi}, \quad \rho_{\rm gw} = \frac{1}{32\pi} \left< \dot{h}_{ij}(t,\vec{x}) \dot{h}^{ij}(t,\vec{x}) \right>,$$

• Cross-power spectral density

$$S_{IJ} = \left\langle \tilde{r}_I^*(f) \tilde{r}_J(f') \right\rangle = \frac{1}{\gamma} \frac{H_0^2}{16\pi^4 f^5} \delta(f - f') \Gamma_{IJ}(f, L_I, L_J, \xi) \,\Omega_{\rm gw}(f)$$

• Overlap reduction function (ORF) is function of f, L_I, L_J, ξ

$$\Gamma_{IJ} = \gamma \sum_{A} \int d\hat{\Omega} \left(e^{2\pi i f L_{I} \left(1 + \hat{\Omega} \cdot \hat{p}_{I} \right)} - 1 \right) \times \left(e^{-2\pi i f L_{J} \left(1 + \hat{\Omega} \cdot \hat{p}_{J} \right)} - 1 \right) F_{I}^{A}(\hat{\Omega}) F_{J}^{A}(\hat{\Omega})$$

• Hellings-Downs correlations for $fL \gg 1$ (short-wavelength approximation)

$$\Gamma_{IJ} = \frac{3}{2} \left(\frac{1 - \cos \xi}{2} \right) \ln \frac{1 - \cos \xi}{2} - \frac{1 - \cos \xi}{8} + \frac{1}{2}$$



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Summary O



PTAs in operation



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Evidence for a GWB in PTA data sets



NANOGrav, 2306.16213 (ApJL); PPTA, 2306.16215 (ApJL)

EPTA+InPTA, 2306.16214 (A&A); CPTA, 2306.16216 (RAA)

Cosmological Interpretation for the PTA Signal

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GWB from SMBHB





(a) plus mode



Hellings-Downs curve

$$\begin{split} \Gamma^{\rm TT}_{ab} &= \; \frac{1}{2} \left[1 + \delta_{ab} + 3 \kappa_{ab} \left(\ln \kappa_{\rm ab} - \frac{1}{6} \right) \right] \\ \kappa_{ab} &\equiv \; (1 - \cos \xi_{ab})/2 \end{split}$$







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NANOGrav, 2306.16213 (ApJL); PPTA, 2306.16215 (ApJL) EPTA+InPTA, 2306.16214 (A&A); CPTA, 2306.16216 (RAA)

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Comparing results from different PTAs



Figure 1. Left: free spectral posteriors for each PTA showing the measured HD-correlated GWB power in several frequency bins under no spectral shape assumption. Each PTA used a different Fourier basis set by their own maximum observing time. The semitransparent gray lines are 100 samples from the joint 2D power-law posterior distribution, showing the spread of power-law models that are consistent with all of the PTA's data. Right: 2D posterior for HD-correlated power-law grarmeters. Contours show 86%, 95%, and 97%, of the posterior mass. The vertical dotted line is at $\gamma = 13/3$.

IPTA, 2309.00693 (ApJ)

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Astro-informed	model from		
Astro-informed	model nom		





• A large eccentricity when GWs begin to dominate the SMBHB evolution.

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see also talks from Shao-Jiang Wang and Qing-Juan Yu

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Astro-informed formation model



• The SGWB from SMBHBs should be detected by LISA, Taiji and TianQin.

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Phase transition





Cosmic string

Domain wall



Scalar-induced GW

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Overview of PTA constraints

TABLE II. Bayes factors (BFs) of the power-law (PL), first order phase transition (PT), domain wall (DW), and cosmic string (CS) models compared to the <u>SMBH</u>Bs model.



• Domain wall model is strongly disfavored.

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Phase transition

$\label{eq:Bubble} \begin{array}{l} \mbox{Bubble collisions} + \mbox{Sound Wave} + \mbox{MHD turbulence} \\ \mbox{see also Shao-Jiang Wang's talk} \end{array}$



TABLE I. The ratio of the vacuum energy density α and critical temperature T_* from five holographical QCD-like models.

Model	QCD matter	Holographic QCD-like model	α	T_* (MeV)
S_1	Heavy static quarks with a zero chemical potential	Hard wall	13.5	122 [133,144]
S_2	Heavy static quarks with a zero chemical potential	Soft wall	4.27	191 [133,144]
S_3	Quarks with a finite chemical potential	Hard wall	32.2	112 [134]
S_4	Quarks with a finite chemical potential	Soft wall	4.56	192 [134]
S_5	Pure gluons	Quenched dynamical holographic QCD	0.611	255 [135]



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log10 a

-9

-8

log₁₀ (f/Hz)

-7

-6

-5

-10

• PTA data prefer pure quark systems under the Jouguet detonations case.

log₁₀ (T_{*}/MeV)

ZCC, Shou-Long Li, Puxun Wu, Hongwei Yu, 2312.01824 (PRD)

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Cosmic String





 The intersection between cosmic strings can lead to reconnections and form loops, which will then decay due to relativistic oscillation and emit gravitational waves.



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/ energy density spe	energy density spectrum of cosmic strings			
10 ⁻⁸ 10 ⁻³ G 10 ⁻¹⁰		$ \begin{array}{c} & & & & \\ & & & & \\ & & & & \\ & & & & $		

10-3

f(Hz)



10-5

• p is the reconnection probability:

10-9

- $\bullet \ p=1 \ {\rm for \ classical \ strings}$
- ${\ensuremath{\, \rm o}\,}\ p<1$ in the string-theory-inspired models

 10^{-7}

10-11

10⁻¹² 10⁻¹¹

GV

 10^{-1}

101

103

 $G\mu = 10^{-14}$

PTA band



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Scalar-induced GW: see Lang Liu's talk





Non-Gaussianity of curvature perturbations

Lang Liu, ZCC, Qing-Guo Huang, 2307.01102 (PRDL)

• Equation of state and sound speed of the early Universe

Lang Liu, **ZCC**, Qing-Guo Huang, 2307.14911 (JCAP) Lang Liu, You Wu, **ZCC**, 2310.16500 (JCAP)

• Speed of GW

ZCC, Jun Li, Lang Liu, Zhu Yi, 2401.09818 (PRDL)

• Distinguish the adiabatic and isocurvature fluctuations

ZCC, Lang Liu, 2402.16781

Sound speed resonance

Jia-Heng Jin, ZCC, Zhu Yi, Zhi-Qiang You, Lang Liu, 2307.08687 (JCAP)



$$\left\langle \tilde{r}_{I}^{*}(f)\tilde{r}_{J}(f')\right\rangle = \frac{1}{\gamma} \frac{H_{0}^{2}}{16\pi^{4}f^{5}} \delta(f-f') \Gamma_{IJ}(f,L_{I},L_{J},\boldsymbol{\xi}) \Omega_{\mathrm{gw}}(f)$$

- PTAs have been opening a new window at nHz frequencies.
- Cosmological implications:
 - Domain wall is strongly disfavored.
 - For PT, PTA data prefer pure quark systems under the Jouguet detonations case.
 - Strings in (super)strings theory are more likely to explain the PTA signal than classical field strings.

Thank you for your attention!