## Constraints on the nonstandard propagating GWs with GWTC-3

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## Cosmological Gravity Theories

- Why modified gravities?
  - Cosmic acceleration
  - Dark matter substitute
  - • •
- Modify weak-field regime (large scales)
- Reduce to GR in strong-field regime by Chameleon/Vainshtein/Symmetron screen mechnisms
- Cosmological tests focus on GW propagation (not generation)



• Even if modification on gravity is a tiny effect, propagation can accumulate the effect because of long distance.



Ezquiaga, Zumalacárregui, Front.Astron.Space Sci. 5 (2018)

- Propagation equation is covariant, i.e. independent of GW sources and background spacetimes (NS, BH, supernova, pulsar, GWB etc.)
- EFT approach [PRD 97 (2018) 10, 104037]

$h_{ij}^{\prime\prime} + \underbrace{(2+\nu)}_{\text{damping}} \mathcal{H} h_i^\prime$	$c_j + \underbrace{c_g^2}_{\text{speed}} k^2 h_{ij} + \underbrace{m}_{\text{dispect}}$	$\sum_{\text{rsion}}^{2} a^{2}h_{ij} =$	oscilltio	$\gamma_{ij}$ ns
gravity theory	ν	$c_{g}^{2} - 1$	$m_g$	Г
GR	0	0	0	0
extra-dim.	$(D-4)\left(1+\frac{1+z}{\mathcal{H}d_{\mathrm{L}}}\right)$	0	0	0
Horndeski	$\alpha_M$	$\alpha_T$	0	0
f(R)	$F'/\mathcal{H}F$	0	0	0
Einstein-aether	0	$c_{\sigma}/\left(1+c_{\sigma}\right)$	0	0
bimetric massive gravity	0	0	$m^2 f_1$	$m^2 f_1$

• Consider  $\Gamma = 0$ 

$$h_{ij}^{\prime\prime} + \underbrace{(2+\nu)}_{\text{damping}} \mathcal{H}h_{ij}^{\prime} + \underbrace{c_g^2}_{\text{speed}} k^2 h_{ij} + \underbrace{m_g^2}_{\text{dispersion}} a^2 h_{ij} = 0$$
(2)

Modified waveform

$$h_{\rm GW} \sim h_{\rm GR} \underbrace{e^{-\frac{1}{2}\int \nu \mathcal{H} d\eta}}_{\text{Affects amplitude}} \underbrace{e^{ik\int \left(c_g^2 - 1 + a^2 m_g^2/k^2\right)^{1/2} d\eta}}_{\text{Affects phase}}$$
(3)

Bonds from GWs

- Bright siren GW170817 (z = 0.008):  $-75.3 \le \nu \le 78.4$  [PRD 97 (2018) 10, 104038]
- GW170817:  $-3 \times 10^{-15} \le c_g 1 \le 7 \times 10^{-16}$  [PRL 119 (2017) 16, 161101]
- GW170104:  $m_q \leq 7.7 \times 10^{-23} \text{eV}$  [PRL 118(22):221101, 2017]

Question: Can we get a tighter constraint on  $\nu$ ?

• Consider  $m_g = \Gamma = 0$  and  $c_g^2 = 1$ 

$$h_{ij}^{\prime\prime} + (2 + \nu)\mathcal{H}h_{ij}^{\prime} + k^2 h_{ij} = 0$$
(4)

Modified luminosity distance

$$d_{\rm GW} = (1+z)^{\nu/2} d_{\rm EM}$$
(5)

$$d_{\rm EM} = \frac{(1+z)}{H_0} \int_0^z \frac{dz'}{\sqrt{\Omega_{\rm m}(1+z')^3 + (1-\Omega_{\rm m})}} \tag{6}$$

ullet GWs measure the luminosity distance  $d_{\rm GW}$  and redshifted masses  $m_1^{\rm det},m_2^{\rm det}$ 

$$m_i = \frac{m_i^{\text{det}}}{1 + z \left( D_{\text{GW}}; H_0, \Omega_{\text{m}} \right)} \tag{7}$$

- Bright siren: infer z with EM counterparts, such as GW170817.
- Dark siren: infer z with galaxy catalogue

## Spectral siren

Even in the absence of electromagnetic observations, GWs alone can probe the expansion rate with the help of population properties, such as

- the peak of the mass distribution;
- the lower/upper mass cut-off;
- redshift distribution.



图: Masses and distance (redshift) distribution from GWTC-3.

## Spectral and bright sirens with GWTC-3 [ApJ 949 (2023) 2, 76]



$$\mathcal{L}(\mathbf{d}|\Lambda) \propto N_{\rm exp}^{N_{\rm obs}} e^{-N_{\rm exp}} \prod_{i=1}^{N_{\rm obs}} \frac{1}{\xi(\Lambda)} \left\langle \frac{\mathcal{R}_{\rm pop}(\theta|\Lambda)}{d_L^2(z)} \right\rangle$$
(8)

• 
$$\mathbf{d} = (d_1, \dots, d_{N_{\mathrm{obs}}})$$
 are  $N_{\mathrm{obs}}$  BBHs

•  $\xi(\Phi)$  quantifies selection biases

$$\xi(\Lambda) = \int P_{\rm det}(\theta) \, p_{\rm pop}(\theta|\Lambda) \, \mathrm{d}\theta \approx \frac{1}{N_{\rm inj}} \sum_{j=1}^{N_{\rm found}} \frac{p_{\rm pop}(\theta_j|\Lambda)}{p_{\rm draw}(\theta_j)}$$

where  $N_{\rm inj}$  is the number of injections,  $N_{\rm found}$  is the number of injections that are detected, and  $p_{\rm draw}\,$  is the probability distribution from which the injections are drawn.

•  $\mathcal{L}(d_i|\theta)$  is single event likelihood.

$$h_{ij}^{\prime\prime} + (2 + \nu)\mathcal{H}h_{ij}^{\prime} + k^2 h_{ij} = 0$$
(9)



- Phenomenological mass models following LVK [ApJ 949 (2023) 2, 76]
- Spectral siren:  $-2.2 \le \nu \le 3.7$  at 90% C.I.
- An oder of magnitude tighter than the bound from bright siren:  $-75.3 \le \nu \le 78.4$